

Turning Weyl's tile argument into a no-go theorem

based on

Velocity polytopes of periodic graphs, arXiv:1109.1963

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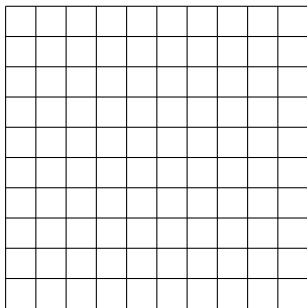
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Perimeter Institute

Experimental Search for Quantum Gravity: the hard facts
Perimeter Institute, Oct 2012

- Proposals for discrete spacetime have problems with Lorentz invariance, but also with isotropy of space.
- Maybe this is only because we haven't found the right models yet?
- I will present a mathematically rigorous no-go theorem which shows that the answer is *No* in a certain sense.
- So far, the theorem only applies to classical point particles.
- Advantages of mathematical rigor:
 - ▶ Requires making explicit definitions
 - ▶ Clear-cut assumptions
 - ▶ You can't argue against it!

Weyl's tile argument

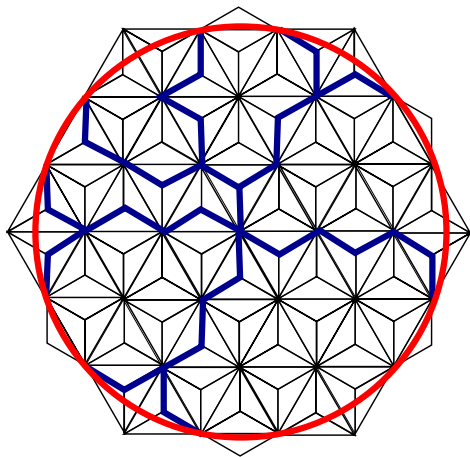


“If a square is built up of miniature tiles, then there are as many tiles along the diagonal as there are along the side; thus the diagonal should be equal in length to the side.”

—Hermann Weyl,
Philosophy of Mathematics and Natural Science,
1927/49

Beyond square lattices

Maybe this is only because a square lattice is too simplistic?



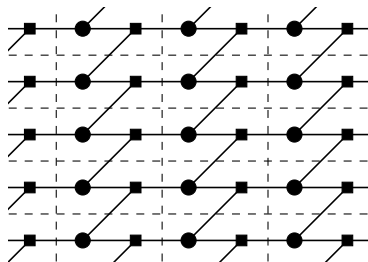
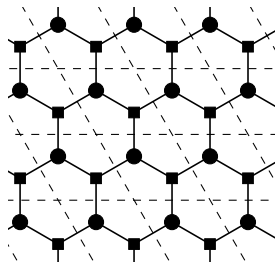
—Giacomo Mauro D’Ariano, *A Quantum-Digital Universe*, 2011

Formalizing the question I: Periodic graphs

Definition

A **periodic graph** is a graph G together with a free action of the discrete translation group \mathbb{Z}^D with finitely many orbits.

→ No ambient space is needed!



Different embeddings of the same periodic graph!

Formalizing the question II: Trajectories

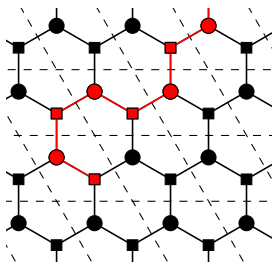
Question

Can physics on a periodic graph be isotropic?

→ This may depend on the kind of physics considered.
In this talk: physics = classical point particles.

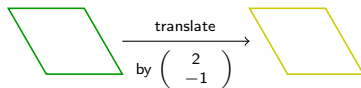
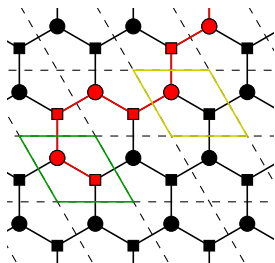
Definition

A **trajectory** is a sequence x_0, x_1, x_2, \dots with $x_t \in G$, indexed by discrete time $t \in \mathbb{N}$, such that every x_t is adjacent to x_{t+1} .



Formalizing the question III: Displacements

Let $C \subseteq G$ be a unit cell. Consider a trajectory $(x_t)_{t \in \mathbb{N}}$ with $x_0 \in C$.



Definition

The **displacement** is the translation $d_t \in \mathbb{Z}^D$ such that $x_t - d_t \in C$.

Formalizing the question IV: Velocities

Definition

The trajectory's **(macroscopic) velocity** $v \in \mathbb{R}^D$ is

$$v = \lim_{t \rightarrow \infty} \frac{d_t}{t},$$

if the limit exists.

Proposition

Velocities do not depend on the choice of unit cell.

→ No notion of distance is required in order to talk about velocities!

Main result

The set of velocities is a subset of \mathbb{R}^D .

→ Isotropy would mean that this set is the unit ball of some Euclidean metric, i.e. an ellipsoid.

Theorem

If G is connected, then the set of velocities is a **convex polytope** in \mathbb{R}^D .

Furthermore, there is a simple algorithm for computing the vertices of that polytope.

→ The continuum limit of a periodic graph, as experienced by a classical point particle, **cannot be isotropic**.